

Stress Intensity Factors for a Crack in Bonded Dissimilar Materials Subjected to Various Stresses

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Abstract The modified complex variable function method with the continuity conditions of the resultant force and displacement function are used to formulate the hypersingular integral equations (HSIE) for an inclined crack and a circular arc crack lies in the upper part of bonded dissimilar materials subjected to various remote stresses. The curve length coordinate method and appropriate quadrature formulas are used to solve numerically the unknown crack opening displacement (COD) function and the traction along the crack as the right hand term of HSIE. The obtained COD is then used to compute the stress intensity factors (SIF), which control the stability behavior of bodies or materials containing cracks or flaws. Numerical results showed the behavior of the nondimensional SIF at the crack tips. It is observed that the nondimensional SIF at the crack tips depend on the various remote stresses, the elastic constants ratio, the crack geometries and the distance between the crack and the boundary.

Keywords Stress Intensity Factor, Bonded Dissimilar Materials, Complex Variable Function, Hypersingular Integral Equation

1. Introduction

A number of papers have been published to analyze the behaviour of stress intensity factors (SIF) at the crack tips subjected to specific remote stress for the crack problems in an infinite plane [1,2], finite plane [3,4], half plane [5,6] or bonded dissimilar materials [7-9]. The body force method with continuous distributions along cracks were used to find the nondimensional SIF of the crack problems in bonded dissimilar materials subjected to various stress

[7]. The Fredholm integral equations with density distributions as undetermined functions were used to calculate the nondimensional SIF for multiple crack problems in bonded dissimilar materials [8]. The logarithmic singular integral equations were used to solve the nondimensional SIF for a circular arc crack lie in the upper part of bonded dissimilar materials [9]. The combinations of Chebyshev polynomials and collocation methods were utilized to solve the nondimensional SIF of a perpendicular crack to the interface of bonded dissimilar materials [10]. The effect of elastic constants ratio to the penny-shaped crack problems in bonded dissimilar materials was determined by reducing the dual integral equations to an infinite system of simultaneous equations [11]. The solidification crack growth patterns in bonded dissimilar materials were predicted by analyzing the nondimensional SIF [12]. The nondimensional SIF of an interface crack in bonded dissimilar materials was evaluated by combining the extended finite element method and a domain independent interaction integral method [13].

The objectives of this paper is to formulate the HSIE and analyze the behavior of nondimensional SIF by using the modified complex variable function method for an inclined and a curved crack lie in the upper part of bonded dissimilar materials subjected to the various remote stresses such as shear stress $\sigma_{x_1} = \sigma_{x_2} = p$, normal stress $\sigma_{y_1} = \sigma_{y_2} = p$, tearing stress $\sigma_{x_1 y_1} = \sigma_{x_2 y_2} = p$ or mixed stress $\sigma_{x_1} = \sigma_{x_2} = \sigma_{y_1} = \sigma_{y_2} = p$.

2. Mathematical Formulation

The complex variable function method plays an important role in solving the cracks problem in plane elasticity [14]. In this method, the stress components

$(\sigma_x, \sigma_y, \sigma_{xy})$, resultant force function $f(X, Y)$ and displacements (u, v) are related to the two complex potentials $\phi(z)$ and $\omega(z)$ as follows

$$\sigma_y - i\sigma_{xy} = \phi'(z) + (z - \bar{z})\overline{\phi''(z)} + \omega'(\bar{z}) \quad (1)$$

$$f = -Y + iX = \phi(z) + (z - \bar{z})\overline{\phi''(z)} + \omega(\bar{z}) \quad (2)$$

$$2G(u + iv) = \kappa\phi(z) - (z - \bar{z})\overline{\phi'(z)} - \omega(\bar{z}) \quad (3)$$

where G is shear modulus of elasticity, $\kappa = (3 - \nu)/(1 + \nu)$ for plane stress, $\kappa = 3 - 4\nu$ for plane strain and ν is Poisson's ratio. The function $\omega(z)$ defined as

$$\omega(z) = z\overline{\phi'(z)} + \overline{\psi(z)} \quad (4)$$

where an analytic function $\overline{\phi}(z)$ defined as $\overline{\phi}(z) = \overline{\phi(\bar{z})}$. The derivative in a specified direction (2) and apply (4), yields

$$\begin{aligned} \frac{d}{dz}\{-Y + iX\} &= \phi'(z) + (z - \bar{z})\overline{\phi''(z)}\frac{d\bar{z}}{dz} + \omega'(\bar{z})\frac{d\bar{z}}{dz} \\ &= \phi'(z) + \overline{\phi'(z)} + \frac{d\bar{z}}{dz}\left[\overline{z\phi''(z)} + \overline{\psi'(z)}\right] = N + iT \end{aligned} \quad (5)$$

where $N + iT$ denotes the normal and tangential traction along the crack segment $\overline{z, z + dz}$. The complex potentials for a crack L in an infinite plane can be expressed by

$$\phi(z) = \frac{1}{2\pi} \int_L \frac{g(t)dt}{t-z}, \quad (6)$$

$$\psi(z) = \frac{1}{2\pi} \int_L \frac{g(t)d\bar{t}}{t-z} + \frac{1}{2\pi} \int_L \frac{\overline{g(t)}dt}{t-z} - \frac{1}{2\pi} \int_L \frac{\bar{t}g(t)dt}{(t-z)^2} \quad (7)$$

where $g(t)$ is crack opening displacement (COD) function defined by

$$g(t) = \frac{2G}{i(\kappa + 1)} \left[(u(t) + iv(t))^+ - (u(t) + iv(t))^- \right] \quad (8)$$

$(u(t) + iv(t))^+$ and $(u(t) + iv(t))^-$ denote the displacements at point t of the upper and lower crack faces, respectively [15].

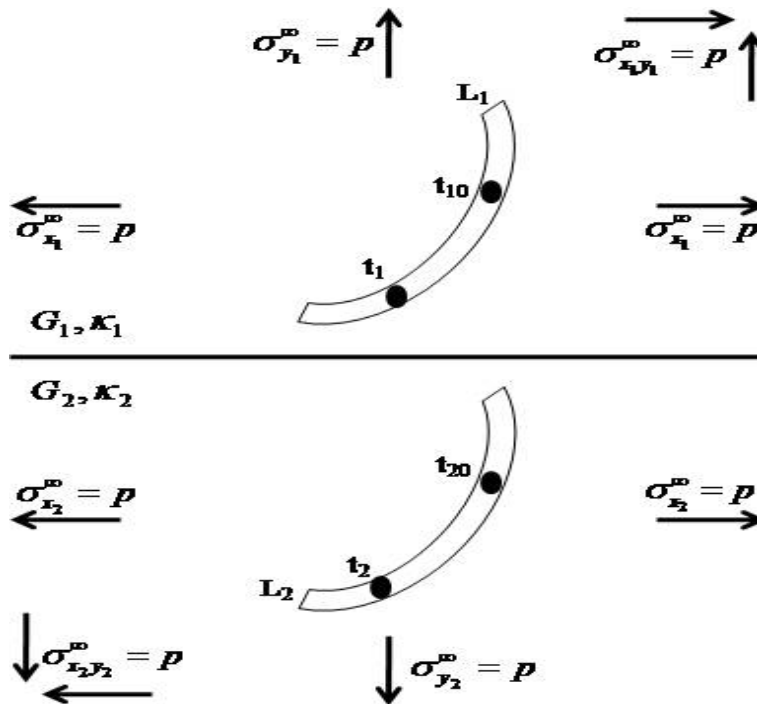


Figure 1. Two cracks L_1 and L_2 in bonded dissimilar materials

Consider two cracks lie in the upper and lower parts of bonded dissimilar materials subjected to various stresses defined in Figure 1. The conditions for remote stresses in bonded dissimilar materials are

$$\varepsilon_{x_1} = \varepsilon_{x_2}, \varepsilon_{y_1} = \varepsilon_{y_2}, \varepsilon_{x_1y_1} = \varepsilon_{x_2y_2}, \varepsilon_{x_1} = \varepsilon_{x_2} = \varepsilon_{y_1} = \varepsilon_{y_2} \quad (9)$$

where subscripts 1 and 2 are the stress component for the upper and lower parts of bonded dissimilar materials, respectively. Those stresses can be expresses as

$$\varepsilon_{x_1} = \frac{1}{E_1}(\sigma_{x_1} - \nu_1 \sigma_{y_1}), \varepsilon_{x_2} = \frac{1}{E_2}(\sigma_{x_2} - \nu_2 \sigma_{y_2}) \quad (10)$$

$$\varepsilon_{y_1} = \frac{1}{E_1}(\sigma_{y_1} - \nu_1 \sigma_{x_1}), \varepsilon_{y_2} = \frac{1}{E_2}(\sigma_{y_2} - \nu_2 \sigma_{x_2}) \quad (11)$$

$$\varepsilon_{x_1y_1} = \frac{1+\nu_1}{E_1} \sigma_{x_1y_1}, \varepsilon_{x_2y_2} = \frac{1+\nu_2}{E_2} \sigma_{x_2y_2} \quad (12)$$

$$\varepsilon_{x_1} = \varepsilon_{y_1} = \frac{1-\nu_1}{E_1} \sigma_{x_1} = \frac{1-\nu_1}{E_1} \sigma_{y_1}, \quad (13)$$

$$\varepsilon_{x_2} = \varepsilon_{y_2} = \frac{1-\nu_2}{E_2} \sigma_{x_2} = \frac{1-\nu_2}{E_2} \sigma_{y_2}$$

where $E_1 = 2G_1(1+\nu_1)$ and $E_2 = 2G_2(1+\nu_2)$ are Young's modulus of elasticity for upper and lower parts of bonded dissimilar materials, respectively. For the shear stress $\sigma_{x_1} = \sigma_{x_2} = p$, normal stress $\sigma_{y_1} = \sigma_{y_2} = p$, tearing stress $\sigma_{x_1y_1} = \sigma_{x_2y_2} = p$ and mixed stress $\sigma_{x_1} = \sigma_{x_2} = \sigma_{y_1} = \sigma_{y_2} = p$ occur when the others stresses does not exist, then (10), (11), (12), (13) reduced to

$$\frac{1}{E_1} \sigma_{x_1} = \frac{1}{E_2} \sigma_{x_2} \quad (14)$$

$$\frac{1}{E_1} \sigma_{y_1} = \frac{1}{E_2} \sigma_{y_2} \quad (15)$$

$$\frac{1+\nu_1}{E_1} \sigma_{x_1y_1} = \frac{1+\nu_2}{E_2} \sigma_{x_2y_2} \quad (16)$$

$$\frac{1-\nu_1}{E_1} \sigma_{x_1} = \frac{1-\nu_2}{E_2} \sigma_{x_2} \text{ or } \frac{1-\nu_1}{E_1} \sigma_{y_1} = \frac{1-\nu_2}{E_2} \sigma_{y_2} \quad (17)$$

respectively.

The modified complex potentials for the crack lies in the upper part of bonded dissimilar materials are defined as

$$\phi_1(z) = \phi_{1p}(z) + \phi_{1c}(z), \psi_1(z) = \psi_{1p}(z) + \psi_{1c}(z) \quad (18)$$

where subscript p and c represent the principal and complementary parts, respectively. Whereas, for the crack lies in the lower part, the complex potentials are represented by $\phi_2(z)$ and $\psi_2(z)$. The continuity conditions for the resultant force (2) and displacement

functions (3) by using (4), yields

$$\left[\phi_1(t) + t \overline{\phi_1'(t)} + \overline{\psi_1(t)} \right]^+ = \left[\phi_2(t) + t \overline{\phi_2'(t)} + \overline{\psi_2(t)} \right]^- \quad (19)$$

$$G_2 \left[\kappa_1 \phi_1(t) - t \overline{\phi_1'(t)} - \overline{\psi_1(t)} \right]^+ = G_1 \left[\kappa_2 \phi_2(t) - t \overline{\phi_2'(t)} - \overline{\psi_2(t)} \right]^- \quad (20)$$

where $t \in L$. Applying (18) into (19) and (20), we can obtained the following complex potentials

$$\phi_{1c}(z) = \frac{G_2 - G_1}{G_1 + \kappa_1 G_2} \left[z \overline{\phi_{1p}'(z)} + \overline{\psi_{1p}(z)} \right], z \in S_1 + L_b \quad (21)$$

$$\psi_{1c}(z) = \frac{\kappa_1 G_2 - \kappa_2 G_1}{G_2 + \kappa_2 G_1} \overline{\phi_{1p}(z)} - \frac{G_2 - G_1}{G_1 + \kappa_1 G_2} \left[z \overline{\phi_{1p}'(z)} + z^2 \overline{\phi_{1p}''(z)} + z \overline{\psi_{1p}'(z)} \right], z \in S_1 + L_b \quad (22)$$

$$\phi_2(z) = \left(1 + \frac{\kappa_1 G_2 - \kappa_2 G_1}{G_2 + \kappa_2 G_1} \right) \phi_{1p}(z), z \in S_2 + L_b \quad (23)$$

$$\psi_2(z) = \left(\frac{G_2 - G_1}{G_1 + \kappa_1 G_2} - \frac{\kappa_1 G_2 - \kappa_2 G_1}{G_2 + \kappa_2 G_1} \right) z \overline{\phi_{1p}'(z)} + \left(1 + \frac{G_2 - G_1}{G_1 + \kappa_1 G_2} \right) \overline{\psi_{1p}(z)}, z \in S_2 + L_b \quad (24)$$

where L_b is boundary of two bonded half materials, S_1 and S_2 are upper and lower parts of bonded dissimilar materials, respectively, and the principle part of complex potentials are referred to an infinite materials.

The HSIE for a crack lies in the upper part of bonded dissimilar materials can be obtained by substituted (18) into (5) and applying (6), (7), (21) and (22), then letting point z approaches t_0 on the crack and changing $d\bar{z}/dz$ into $d\bar{t}_0/dt_0$, yields

$$N(t_0) + iT(t_0) = \frac{1}{\pi} h p \int_L \frac{g(t) dt}{(t-t_0)^2} + \frac{1}{2\pi} \int_L A_1(t, t_0) g(t) dt + \frac{1}{2\pi} \int_L A_2(t, t_0) \overline{g(t)} dt \quad (25)$$

where

$$A_1(t, t_0) = B_1(t, t_0) + \frac{G_2 - G_1}{G_1 + \kappa_1 G_2} \left[\overline{B_3(t, t_0)} - \frac{d\bar{t}_0}{dt_0} (B_5(t, t_0) + \overline{B_4(t, t_0)}) + \frac{d\bar{t}}{dt} (B_4(t, t_0) + \overline{B_4(t, t_0)}) - \frac{d\bar{t} d\bar{t}_0}{dt dt_0} B_6(t, t_0) \right] + \frac{\kappa_1 G_2 - \kappa_2 G_1}{G_2 + \kappa_2 G_1} \frac{d\bar{t}_0}{dt_0} \overline{B_4(t, t_0)}$$

$$A_2(t, t_0) = B_2(t, t_0) + \frac{G_2 - G_1}{G_1 + \kappa_1 G_2} \left[\overline{B_4(t, t_0)} + \overline{B_4(t, t_0)} - \frac{d\bar{t}_0}{dt_0} B_6(t, t_0) + \frac{d\bar{t}}{dt} B_3(t, t_0) \right]$$

and

$$\begin{aligned}
 B_1(t, t_0) &= \frac{1}{(t-t_0)^2} \left[\frac{(t-t_0)^2}{(\bar{t}-\bar{t}_0)^2} \frac{d\bar{t}}{dt} \frac{d\bar{t}_0}{dt_0} - 1 \right] \\
 B_2(t, t_0) &= \frac{t-t_0}{(\bar{t}-\bar{t}_0)^3} \left[\frac{(\bar{t}-\bar{t}_0)}{(t-t_0)} \left(\frac{d\bar{t}}{dt} + \frac{d\bar{t}_0}{dt_0} \right) - 2 \frac{d\bar{t}}{dt} \frac{d\bar{t}_0}{dt_0} \right] \\
 B_3(t, t_0) &= \frac{1}{(\bar{t}-\bar{t}_0)^2} + \frac{2(t_0-t)}{(\bar{t}-\bar{t}_0)^3} \\
 B_4(t, t_0) &= \frac{1}{(\bar{t}-\bar{t}_0)^2} \\
 B_5(t, t_0) &= \frac{2(3\bar{t}_0-2t_0-\bar{t})}{(t-\bar{t}_0)^3} + \frac{6(\bar{t}_0-\bar{t})(\bar{t}_0-t_0)}{(t-\bar{t}_0)^4} \\
 B_6(t, t_0) &= \frac{1}{(t-\bar{t}_0)^2} + \frac{2(\bar{t}_0-t_0)}{(t-\bar{t}_0)^3}.
 \end{aligned}$$

Note that if $G_2 = 0$ and $G_I = G_2$, then the HSIE for a crack in bonded dissimilar (25) reduces to the HSIE for a crack in the half plane [5] and in an infinite plane elasticity [2], respectively.

It is well known that to solve the HSIE (25), the curved length coordinate method is used which transform the COD function $g(t)$ to the square root singularity $\sqrt{a^2-s^2}$ as follows [5,15]

$$H(s) = \frac{g(t)}{\sqrt{a^2-s^2}} \tag{26}$$

with the following quadrature formulas [16]

$$\frac{1}{\pi} \int_{-a}^a \frac{\sqrt{a^2-s^2} H(s) ds}{(s-s_0)^2} = \sum_{j=1}^{M+1} W_j(s_0) H(s_j) \tag{27}$$

$$\frac{1}{\pi} \int_{-a}^a \sqrt{a^2-s^2} H(s) ds = \frac{1}{M+2} \sum_{j=1}^{M+1} (a^2-s_0^2) H(s_j) \tag{28}$$

where $H(s)$ is a given function, $M \in \mathbb{Z}^+$,

$$s_j = a \cos\left(\frac{j\pi}{M+2}\right), j = 1, 2, \dots, M+1$$

and

$$W_j(s_0) = -\frac{2}{M+2} \sum_{n=0}^M (n+1) \sin\left(\frac{j\pi}{M+2}\right) \sin\left(\frac{(n+1)j\pi}{M+2}\right) U_n\left(\frac{s_{j0}}{a}\right)$$

and the observation points

$$s_0 = s_{0,k} = a \cos\left(\frac{k\pi}{M+2}\right), k = 1, 2, \dots, M+1.$$

Here $U_n(t)$ is a Chebyshev polynomial of the second kind, define by

$$U_n(t) = \frac{\sin((n+1)\theta)}{\sin\theta}, \text{ where } t = \cos\theta. \tag{29}$$

The normal and tangential components $N + iT$ of the HSIE for the shear stress $\sigma_{x_1} = \sigma_{x_2} = p$, normal stress $\sigma_{y_1} = \sigma_{y_2} = p$, tearing stress $\sigma_{x_1y_1} = \sigma_{x_2y_2} = p$ and mixed stress $\sigma_{x_1} = \sigma_{x_2} = \sigma_{y_1} = \sigma_{y_2} = p$ with an angle of the crack α are defined as follows

$$N + iT = -p \sin^2 \alpha - ip \sin \alpha \cos \alpha \tag{30}$$

$$N + iT = -p \cos^2 \alpha + ip \sin \alpha \cos \alpha \tag{31}$$

$$N + iT = 2p \sin \alpha \cos \alpha + ip (\cos^2 \alpha - \sin^2 \alpha) \tag{32}$$

$$N + iT = -p + i0. \tag{33}$$

2. Results and Discussions

Stress intensity factor (SIF) at crack tips A_j is defined as

$$(K_1 - iK_2)_{A_j} = \sqrt{2\pi} \lim_{t \rightarrow t_{A_j}} \sqrt{|t-t_{A_j}|} g'(t), j = 1, 2. \tag{34}$$

Consider an inclined crack lie in the upper part of bonded dissimilar materials subjected to various stresses as defined in Figure 2.

Table 1 displays the nondimensional SIF versus $h/2R$ at the crack tip subjected to tearing stress $\sigma_{x_1y_1} = \sigma_{x_2y_2} = p$ when with $\alpha = 0^\circ$ different values of elastic constant ratio G_2/G_1 where R and h are defined as in Figure 2. Our results are totally agrees with those of Isida and Noguchi [7]. It is observed that the nondimensional SIF at F_{1A1} and F_{2A1} are equals to negative of F_{1A2} and F_{2A2} , respectively.

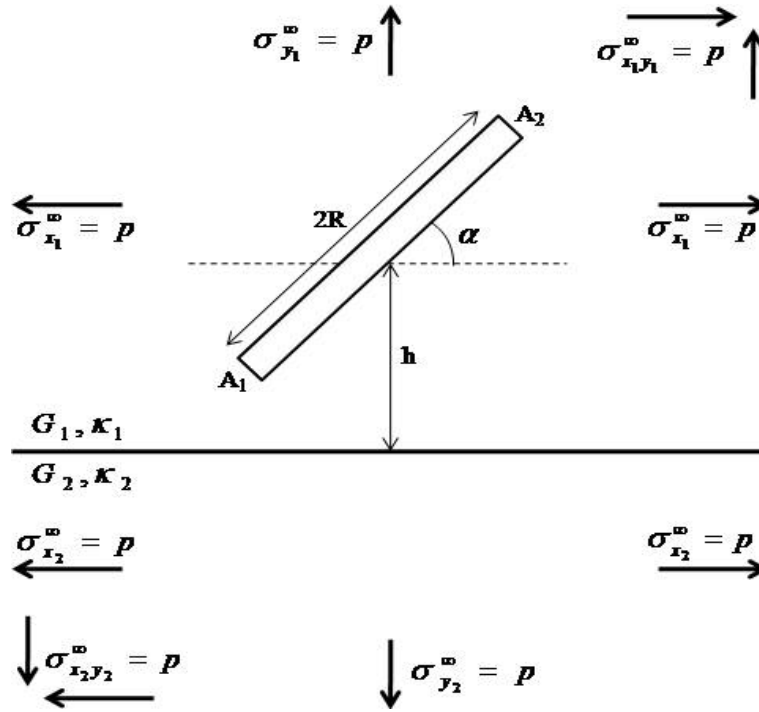


Figure 2. An inclined crack in bonded dissimilar materials

Table 1. Nondimensional SIF for an inclined crack when $\alpha = 0^\circ$ subjected to $\sigma_{x_1y_1} = \sigma_{x_2y_2} = p$

G_2/G_1	SIF	$h/2R$								
		0.1	0.2	0.3	0.4	0.5	0.6	0.7	0.8	0.9
0.0	F_{1A1}^*	0.6151	0.3876	0.2639	0.1854	0.1331	0.0973	0.0724	0.0548	0.0422
	F_{1A1}^{**}	0.6150	0.3880	0.2640	0.1850	0.1330	0.0970	0.0720	0.0550	0.0420
	F_{2A1}^*	1.2077	1.1329	1.1126	1.0993	1.0874	1.0762	1.0661	1.0573	1.0497
	F_{2A1}^{**}	1.2080	1.1330	1.1130	1.0990	1.0870	1.0760	1.0660	1.0570	1.0500
0.5	F_{1A1}^*	0.0878	0.0658	0.0516	0.0398	0.0303	0.0230	0.0175	0.0135	0.0105
	F_{1A1}^{**}	0.0880	0.0660	0.0520	0.0400	0.0300	0.0230	0.0170	0.0130	0.0100
	F_{2A1}^*	1.0864	1.0512	1.0372	1.0303	1.0257	1.0221	1.0190	1.0164	1.0142
	F_{2A1}^{**}	1.0860	1.0510	1.0370	1.0300	1.0260	1.0220	1.0190	1.0160	1.0140
2.0	F_{1A1}^*	-0.0696	-0.0514	-0.0403	-0.0316	-0.0245	-0.0188	-0.0145	-0.0112	-0.0087
	F_{1A1}^{**}	-0.0700	-0.0510	-0.0400	-0.0320	-0.0240	-0.0190	-0.0140	-0.0110	-0.0090
	F_{2A1}^*	0.9176	0.9456	0.9602	0.9684	0.9738	0.9778	0.9811	0.9838	0.9861
	F_{2A1}^{**}	0.9180	0.9460	0.9600	0.9680	0.9740	0.9780	0.9810	0.9840	0.9860

*Present study

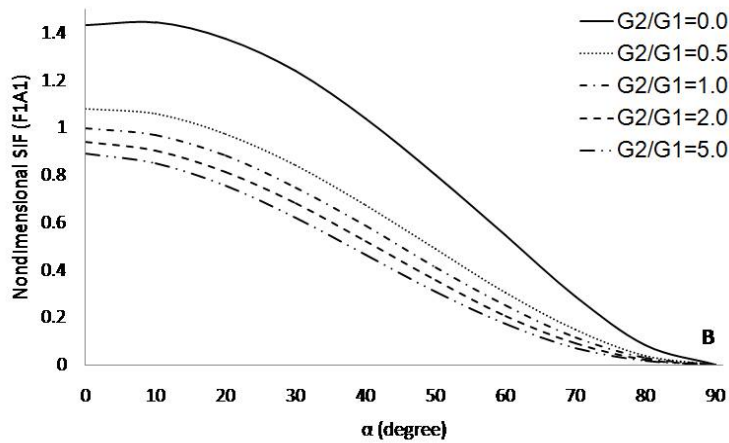
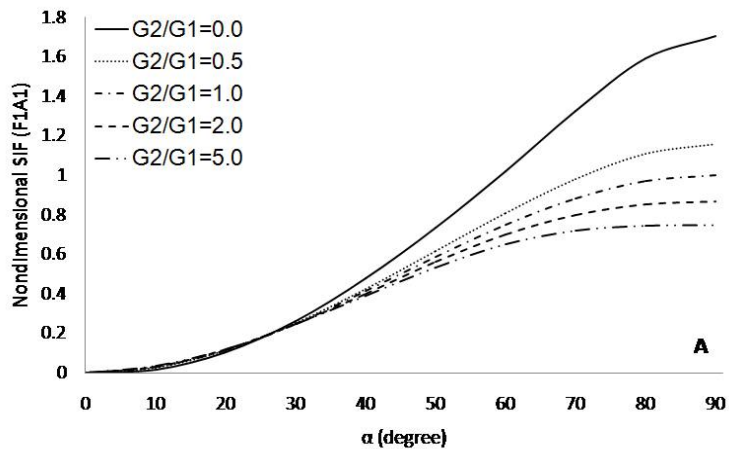
**Isida and Noguchi [7]

The Figures 3 and 4 present the nondimensional SIF for an inclined crack lies in upper part of bonded dissimilar materials subjected to various stresses at crack tips A_1 and A_2 , respectively, when $R/h = 0.9$ and α varies (Figure 2). It is found that as α increases the nondimensional SIF increases subjected to shear stress $\sigma_{x_1} = \sigma_{x_2} = p$ and decreases subjected to normal stress $\sigma_{y_1} = \sigma_{y_2} = p$ at crack tips A_1 and A_2 . For tearing stress $\sigma_{x_1y_1} = \sigma_{x_2y_2} = p$ the nondimensional SIF increases when $\alpha > 50^\circ$ at crack tips A_1 and A_2 . Whereas as α increases the nondimensional SIF at crack tips A_1 and A_2 does not show any significant differences for $G_2/G_1 > 0.0$ subjected to mixed stress $\sigma_{x_1} = \sigma_{x_2} = \sigma_{y_1} = \sigma_{y_2} = p$. The nondimensional SIF decreases as G_2/G_1 increases subjected to shear stress $\sigma_{x_1} = \sigma_{x_2} = p$, normal stress $\sigma_{y_1} = \sigma_{y_2} = p$ and mixed stress $\sigma_{x_1} = \sigma_{x_2} = \sigma_{y_1} = \sigma_{y_2} = p$ at crack tips A_1 and A_2 . However for tearing stress $\sigma_{x_1y_1} = \sigma_{x_2y_2} = p$ the

nondimensional SIF increases as G_2/G_1 increases at crack tips A_1 and A_2 .

Consider a circular arc crack lie in the upper part of bonded dissimilar materials subjected to various stresses as defined in Figure 5.

Figure 6 shows the nondimensional SIF for a circular arc crack lies in the upper part of bonded dissimilar materials subjected to various stresses at crack tip A_1 when $h = 0.5R$ and α varies as defined in Figure 5. It is observed that the nondimensional SIF at crack A_1 equal to the SIF A_2 subjected to shear stress $\sigma_{x_1} = \sigma_{x_2} = p$, normal stress $\sigma_{y_1} = \sigma_{y_2} = p$ and mixed stress $\sigma_{x_1} = \sigma_{x_2} = \sigma_{y_1} = \sigma_{y_2} = p$. Whereas for tearing stress $\sigma_{x_1y_1} = \sigma_{x_2y_2} = p$ the nondimensional SIF at crack A_1 equal to the negative SIF at A_2 . The nondimensional SIF increases subjected to shear stress $\sigma_{x_1} = \sigma_{x_2} = p$ and tearing stress $\sigma_{x_1y_1} = \sigma_{x_2y_2} = p$, however SIF decreases subjected to normal stress $\sigma_{y_1} = \sigma_{y_2} = p$ and mixed stress $\sigma_{x_1} = \sigma_{x_2} = \sigma_{y_1} = \sigma_{y_2} = p$ as α increases at crack tip A_1 .



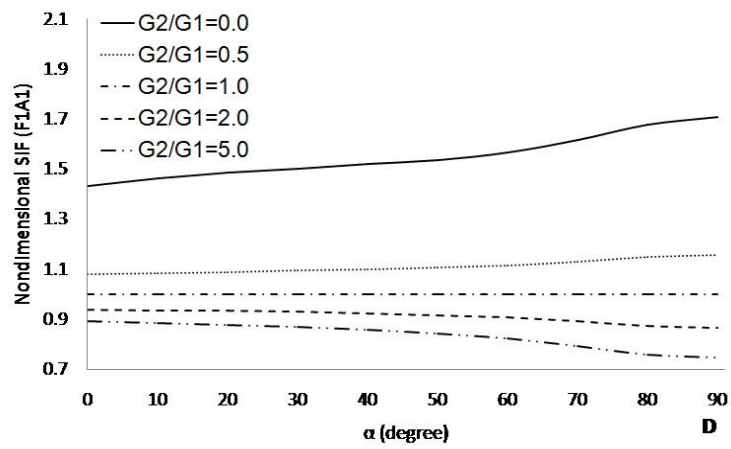
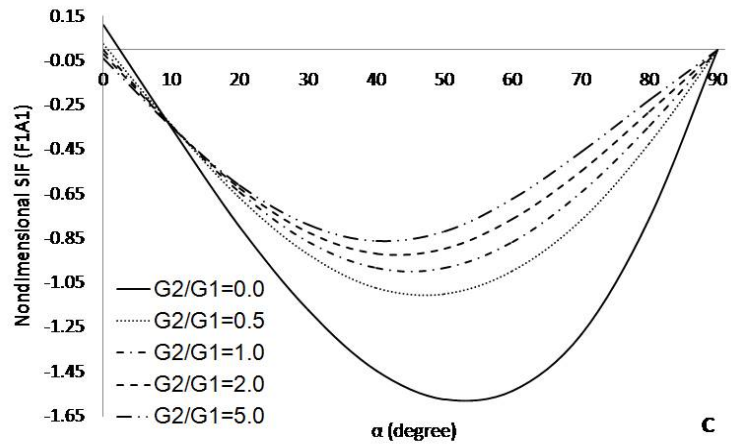
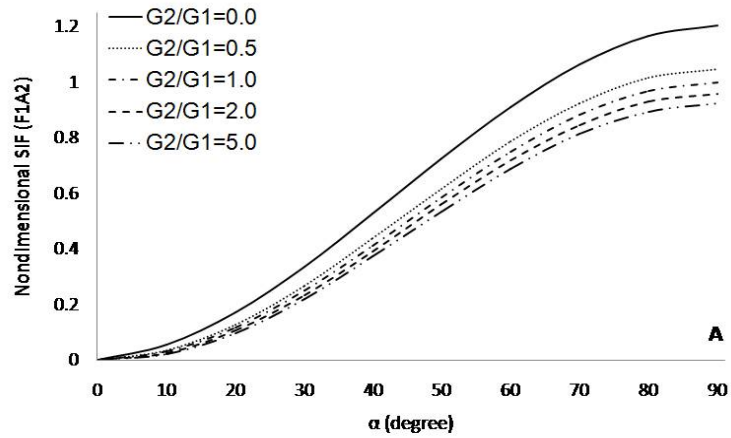


Figure 3. Nondimensional SIF versus α when $R/h = 0.9$ at crack tip A_1 subject to: Part A Shear stress, Part B Normal stress, Part C Tearing Stress and Part D Mixed stress



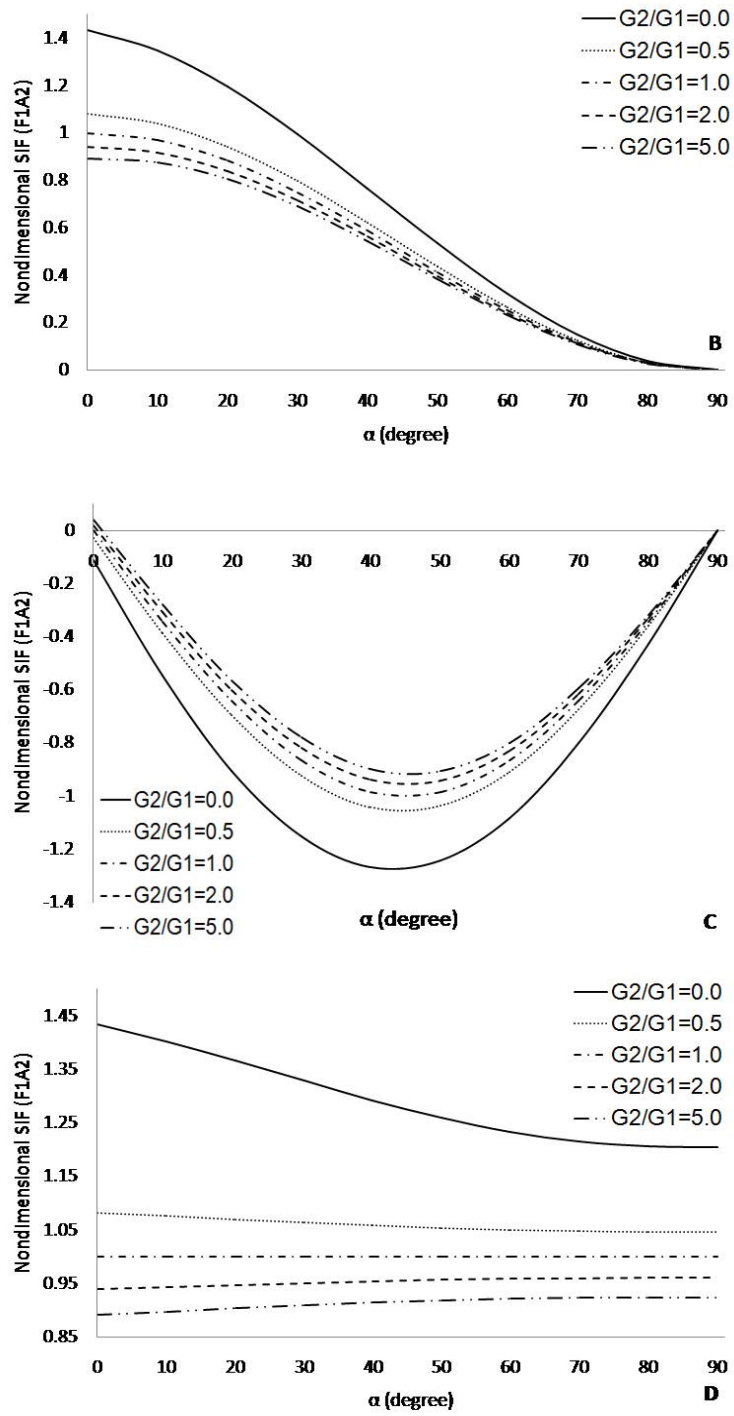


Figure 4. Nondimensional SIF versus α when $R/h = 0.9$ at crack tip A_2 subject to: Part A Shear stress, Part B Normal stress, Part C Tearing Stress and Part D Mixed stress

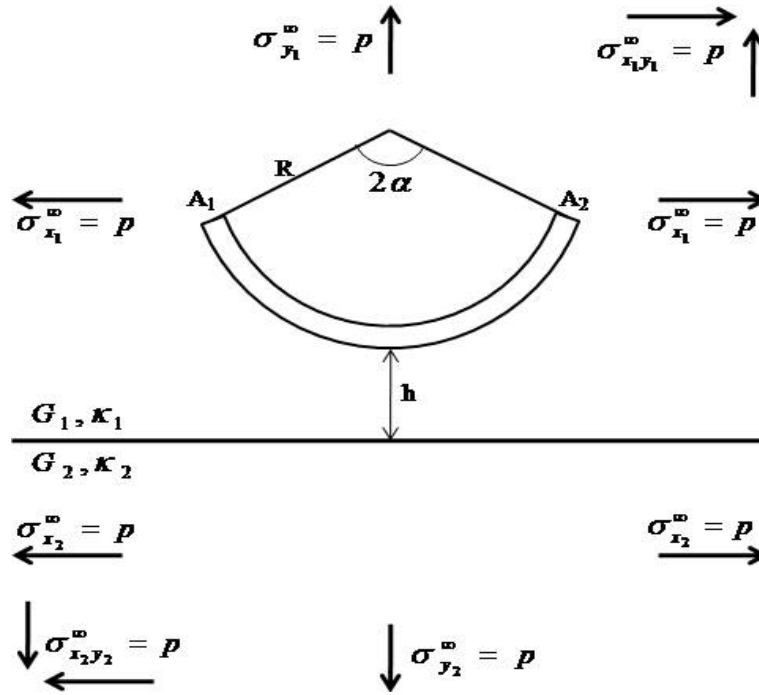
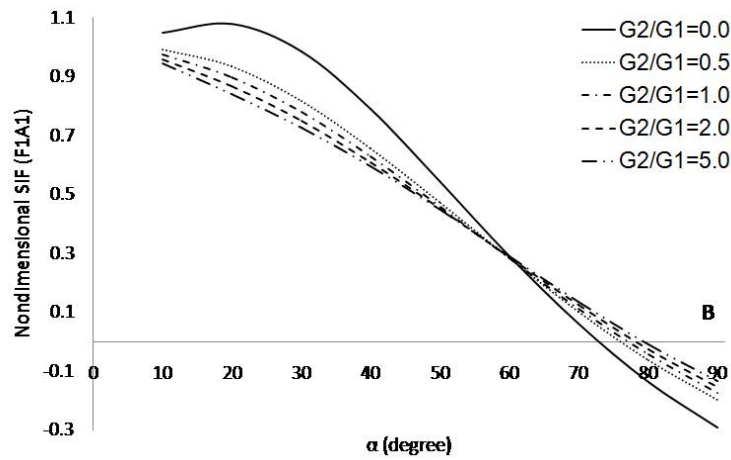
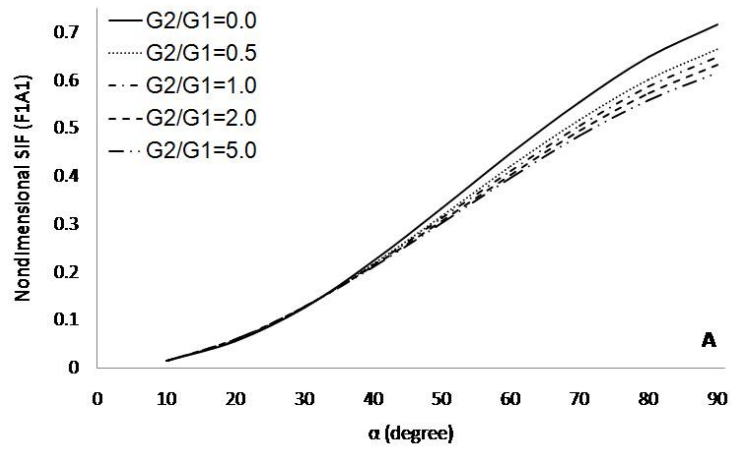


Figure 5. A circular arc crack in bonded dissimilar materials



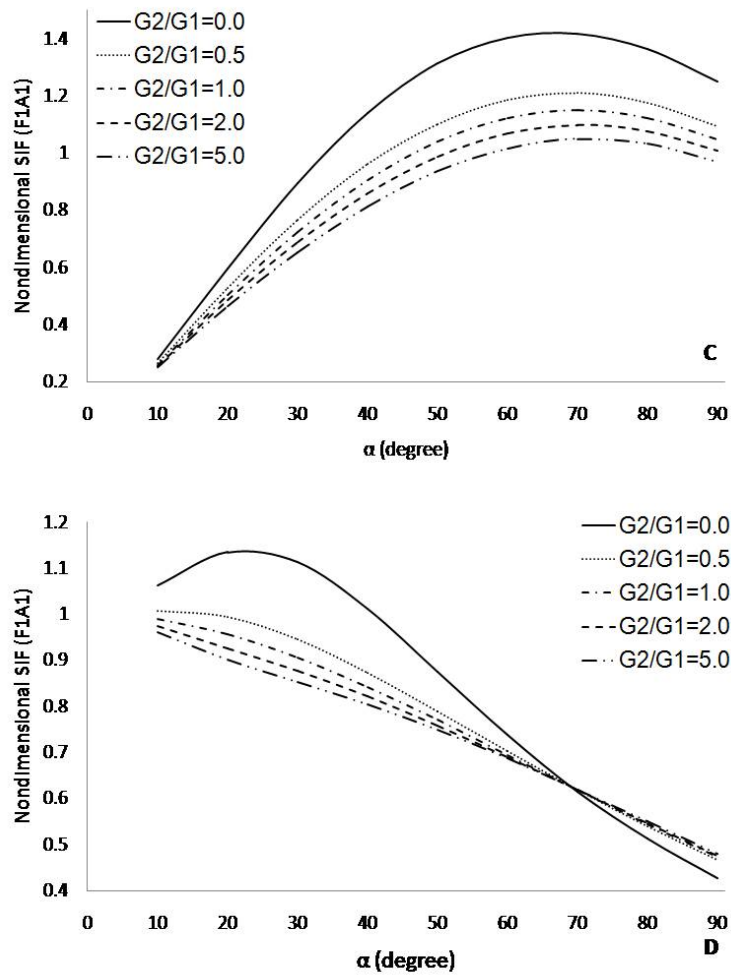


Figure 6. SIF for a circular arc crack when $h = 0.5R$ and α varies at crack tip A_1 subject to: Part A Shear stress, Part B Normal stress, Part C Tearing Stress and Part D Mixed stress

As G_2/G_1 increases the nondimensional SIF decreases subjected to shear stress $\sigma_{x_1} = \sigma_{x_2} = p$ and tearing stress $\sigma_{x_1 y_1} = \sigma_{x_2 y_2} = p$ at crack tip A_1 . Whereas for normal stress $\sigma_{y_1} = \sigma_{y_2} = p$ the nondimensional SIF decreases when $\alpha < 60^\circ$ and for mixed stress $\sigma_{x_1} = \sigma_{x_2} = \sigma_{y_1} = \sigma_{y_2} = p$ when $\alpha < 70^\circ$ as G_2/G_1 increases at crack tip A_1 .

From the numerical results we conclude that the nondimensional SIF for a crack lies in the upper part of bonded dissimilar materials depends on the various remote stresses, the elastic constants ratio, the crack geometries and the distance between the crack and the boundary. The closer the crack to the boundary the smaller the nondimensional SIF which interprets the weaker of the materials. However the nondimensional SIF is altered for different stresses, elastic constants ratio or crack geometries.

4. Conclusions

In this paper, the HSIE for an inclined crack and a circular arc crack lies in the upper part of bonded dissimilar materials subjected to various remote stresses such as shear stress $\sigma_{x_1} = \sigma_{x_2} = p$, normal stress $\sigma_{y_1} = \sigma_{y_2} = p$, tearing stress $\sigma_{x_1 y_1} = \sigma_{x_2 y_2} = p$ or mixed stress $\sigma_{x_1} = \sigma_{x_2} = \sigma_{y_1} = \sigma_{y_2} = p$ are formulated by using the modified complex variable function method with the crack opening displacement function as the unknown function.

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