Numerical Solutions of MHD Viscous Flow of Newtonian Fluids due to a Shrinking Sheet by SOR Iterative Procedure

Mohammad Shafique¹,* , Fatima Abbas²

¹Ex-AP, Department of Mathematics, Gomal University, D I Khan, Pakistan
²Department of Mathematics, Gomal University, D I Khan, Pakistan

Abstract  The problem of Magneto Hydrodynamic viscous flow due to a shrinking sheet of Newtonian fluids has been solved numerically by using SOR Iterative Procedure. The similarity transformations have been used to reduce the highly nonlinear partial differential equations of motion to ordinary differential equations. The results have been calculated on three different grid sizes to check the accuracy of the results. The numerical results for Newtonian fluids are found in good agreement with those obtained by the previous results.

Keywords  Newtonian Fluids, Shrinking Sheet and SOR Iterative Procedure

AMS Subject Classification: 76D99, 76M20, 65N22

1. Introduction


In this research, the numerical solutions of MHD viscous flow due to a shrinking sheet for Newtonian fluid have been discussed. In order to find the numerical solution of the problem, the Navier Stokes equations are reduced to ordinary differential equations by using similarity transformations [13]. This system is solved numerically by using SOR Iterative Procedure with Simpson (1/3) Rule. The calculations have been carried out using three different grid sizes to check the accuracy of the results. The present numerical results have also been compared with the previous results in a particular case and found in good agreement. The numerical results have been discussed in both tabular as well as graphically.

2. Mathematical Analysis

The continuity equation and the Navier-Stokes equations for incompressible fluid in the presence of body forces are given by

\[ \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{V}) = 0, \]  
(1)

\[ \rho \left( \frac{\partial \mathbf{V}}{\partial t} + (\mathbf{V} \cdot \nabla) \mathbf{V} \right) = -\nabla p + \mu \nabla^2 \mathbf{V}. \]  
(2)

Where \( \rho \) and \( \mathbf{V} \) are respectively, the density and the velocity vector of the fluid.

The assumptions are made for the problem under consideration. The fluid flow is steady, laminar and incompressible. The fluid is electrically conducting in the presence of a magnetic field of strength \( B_0 \). The
electromagnetic body force is given as \( f = -\sigma B_0^2 (u, v, 0) \).

The fluid flow in the frame of three dimensional Cartesian coordinate systems and \( \mathbf{V} = \{ u(x, y, z), v(x, y, z), w(x, y, z) \} \) is the velocity field of flow.

Under the above assumptions, the equations (1) and (2) become

\[
\frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} = -\frac{1}{\rho} \frac{\partial p}{\partial x} + v \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} \right) - \frac{\sigma B_0^2}{\rho} u, \tag{3}
\]

\[
\frac{\partial v}{\partial x} + u \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} = -\frac{1}{\rho} \frac{\partial p}{\partial y} + v \left( \frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} + \frac{\partial^2 v}{\partial z^2} \right) - \frac{\sigma B_0^2}{\rho} v, \tag{4}
\]

\[
\frac{\partial w}{\partial x} + u \frac{\partial w}{\partial y} + w \frac{\partial w}{\partial z} = -\frac{1}{\rho} \frac{\partial p}{\partial z} + v \left( \frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 w}{\partial y^2} + \frac{\partial^2 w}{\partial z^2} \right) - \frac{\sigma B_0^2}{\rho} w \tag{5}
\]

Where \( p \) is the pressure and \( \mu \) is Kinematics Viscosity and \( \mu \) denotes viscosity coefficient. The boundary conditions are

\[
u = ax, v = -(m-1)y, w = -W \text{ at } y = 0, \tag{7}
\]

Where \( a > 0 \) is the shrinking constant and \( W \) is the suction velocity. When \( m=1 \), the sheet shrinks in the \( x \)-direction and when \( m=2 \), the sheet shrinks axisymmetrically.

In order to solve equations (3) to (6), the similarity transformations of [13] are given:

\[
u = xaf'(\eta), v = y(m-1)f'(\eta), w = -m\sqrt{av} f(\eta), \tag{8}
\]

where \( \eta = \sqrt{\frac{u}{v}} \) is a dimensionless variable. The continuity equation (3) satisfied and the equation (6) can be integrated to give

\[
\frac{p}{\rho} = \frac{\partial w}{\partial z} - \frac{w^2}{2} \tag{4}
\]

The resulting partial differential equations by using (8) we obtain:

\[
f'' - M^2 f' - f'^2 + mff'' = 0, \tag{9}
\]

with the boundary conditions:

\[
f = S, f' = -1 \text{ at } \eta = 0, \]

\[
f' \to 0 \text{ as } \eta \to \infty. \tag{10}
\]

Where \( S = \frac{W}{m\sqrt{av}} \) and \( M^2 = \frac{\sigma B_0^2}{\rho \mu} \). The prime denotes the differentiation with respect to \( \eta \).

3. Finite Difference Equations

For numerical purpose, we rewrite the equations (9) and (10) by putting

\[
P = f' \tag{11}
\]

The equation (9) as follows:

\[
P'^2 - M^2 P - P^2 + mff' = 0, \tag{12}
\]

and the boundary conditions (10) takes the form:

\[
f = S, P = -1 \text{ at } \eta = 0, \]

\[
P \to 0 \text{ as } \eta \to \infty. \tag{13}
\]

Now if we approximate equation (12) by central difference approximations at a typical point \( \eta = \eta_n \) of the interval \([0, \infty)\), we obtained:

\[
P_{n+1} - 2P_n + P_{n-1} - M^2 P_n + mP_n\left( P_{n+1} - P_{n-1} \right) - P_n^2 = 0, \tag{14}
\]

Where \( h \) denotes a grid size and the symbols used denote as \( f_n = f(\eta_n), P_n = P(\eta_n) \), at a typical point of the interval \([0, \infty)\) replaced by \([0, \beta)\), where \( \beta \) is sufficiently large.

4. Computational Procedure

We now integrate numerically equation (11) and solve the system of finite difference equation (14) at each required grid point of the interval \([0, \infty)\). The equation (11) is integrated by using the Simpson’s (1/3) rule [15] as given below at the gird point \( \eta = \eta_n \)

\[
f_{n+1} = f_n + \frac{h}{3} \left( P_{n+1} + 4P_n + P_{n-1} \right). \tag{15}
\]

Whereas the system of finite-difference equations (14) are solved by using S.O.R. iterative procedure [15]

\[
P_n = \frac{\omega}{(4 + 2h^2 M^2 + 2h^2 P_n)} \left[ 2(P_{n+1} + P_{n-1}) + mP_n(P_{n+1} - P_{n-1}) \right] + (1-\omega) P_n \tag{16}
\]

Subject to the appropriate boundary conditions (13), the computation has been checked for different of the relaxation parameter \( \omega \) between 1 and 2. The optimum value of the relaxation parameter for the problem under consideration is 1.5.

The SOR iterative procedure is terminated when the following criterion is satisfied:

\[
\max_{i=1}^n \left| U_i^{n+1} - U_i^n \right| < 10^{-6} \tag{17}
\]

where \( n \) denotes the number of iterations and \( U \) stands for each of \( P \) and \( f \). The stability of SOR iterative procedure and accuracy of the solution has been checked for step sizes \( h=0.01, 0.005 \) and \( 0.0025 \) and it is not necessary to discuss because this study has been focus over numerical solution of the problem however accuracy is excellent.
5. Numerical Results and Discussions

The numerical computation has been performed to study the effect of the flow parameters namely S and M. The accuracy of the results is checked by comparing them on different grid sizes. The results for the non-dimensional velocity components $f$, $f'$ are shown in the tables 1-3, for each of grid sizes mentioned above. When $m=1$, the sheet shrinks in the x-direction and when $m=2$, the sheet shrinks axisymmetrically. The results compare very well. The comparison of the present results for $f'(0)$ (the skin friction) with the previous results by [13, 14] is given in table 4.

Graphically, the curves for the function $f$, $f'$ have been plotted and are displayed in figures 1 to 8 for several values of the parameters $m$, $S$ and $M$. When $m=1$, the two dimensional case, figures 1-2, with $M=2$ and for different values of $S$, and figures 3-4, with $S=1$ and for different values of $S$. When $m=2$, the axisymmetric case, figures 5-6, with $M=2$ and for different values of $S$, and figures 7-8, with $S=1$ and for different values of $S$.

It can be observed that the function $f'$ increases with the increasing values of $S$ at a fixed value of $M$. The function $f$ decreases initially and then becomes uniform for all values of $S$ and $M$.

Consequently, the effects of different parameters are observed on the similarity function and the velocity $f'$. From this analysis, the boundary layer becomes thinner for larger $S$. The increasing values of $M$ show stronger effect on both the velocity functions. Also, the present results for $f''(0)$ (the skin friction) have good agreement with the previous results.

### Table 1. Numerical Results of $f$ and $f'$ using SOR Iterative Procedure Simpson’s Rule for specified values of parameters $M=1.0$ and $S=0.1, 2.0$

<table>
<thead>
<tr>
<th>$h$</th>
<th>$\eta$</th>
<th>$M=1.0$ and $S=0.1$</th>
<th>$M=1.0$ and $S=2.0$</th>
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<tr>
<td></td>
<td></td>
<td>$f(\eta)$</td>
<td>$f'(\eta)$</td>
</tr>
<tr>
<td>0.01</td>
<td>0.0</td>
<td>0.100000</td>
<td>-1.000000</td>
</tr>
<tr>
<td></td>
<td>1.0</td>
<td>-0.827905</td>
<td>-0.858721</td>
</tr>
<tr>
<td></td>
<td>5.0</td>
<td>-3.228485</td>
<td>0.000000</td>
</tr>
<tr>
<td>0.005</td>
<td>0.0</td>
<td>0.100000</td>
<td>-1.000000</td>
</tr>
<tr>
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<td>1.0</td>
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<td>-0.854891</td>
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<td></td>
<td>5.0</td>
<td>-3.228485</td>
<td>0.000000</td>
</tr>
<tr>
<td>0.0025</td>
<td>0.0</td>
<td>0.100000</td>
<td>-1.000000</td>
</tr>
<tr>
<td></td>
<td>1.0</td>
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</tr>
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### Table 2. Numerical Results of $f$ and $f'$ using SOR Iterative Procedure Simpson’s Rule for specified values of parameters $M=2.0$ and $S=0.5, 1.0$

<table>
<thead>
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<th>$M=2.0$ and $S=1.0$</th>
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<td></td>
<td></td>
<td>$f(\eta)$</td>
<td>$f'(\eta)$</td>
</tr>
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<td>0.500000</td>
<td>-1.000000</td>
</tr>
<tr>
<td></td>
<td>1.0</td>
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<td>-0.135336</td>
</tr>
<tr>
<td></td>
<td>5.0</td>
<td>0.000050</td>
<td>0.000000</td>
</tr>
<tr>
<td>0.005</td>
<td>0.0</td>
<td>0.500000</td>
<td>-1.000000</td>
</tr>
<tr>
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<td>1.0</td>
<td>0.067665</td>
<td>-0.135336</td>
</tr>
<tr>
<td></td>
<td>5.0</td>
<td>0.000050</td>
<td>0.000000</td>
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<tr>
<td>0.0025</td>
<td>0.0</td>
<td>0.500000</td>
<td>-1.000000</td>
</tr>
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<td>1.0</td>
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</table>
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Table 3. Numerical Results of $f$ and $f'$ using SOR Iterative Procedure Simpson’s Rule for specified values of parameters $M=3.0$ and $S=1.0, 1.5$

<table>
<thead>
<tr>
<th>$h$</th>
<th>$\eta$</th>
<th>$m=1$</th>
<th>$m=2$</th>
<th>$m=1$</th>
<th>$m=2$</th>
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<td>1.500000</td>
</tr>
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<td>0.713634</td>
<td>0.713634</td>
<td>1.234863</td>
<td>1.234863</td>
</tr>
<tr>
<td>2.0</td>
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<td>1.228150</td>
<td>1.228150</td>
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<tr>
<td>3.0</td>
<td>0.703472</td>
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<tr>
<td>4.0</td>
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<td>0.703461</td>
<td>0.703461</td>
<td>1.227977</td>
<td>1.227977</td>
</tr>
<tr>
<td>5.0</td>
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<td>0.703461</td>
<td>0.703461</td>
<td>1.227977</td>
<td>1.227977</td>
</tr>
</tbody>
</table>

Table 4. Comparison of the Numerical Value of $f''(0)$ for $M=2.0$

$$
\begin{array}{|c|c|c|c|c|c|}
\hline
S & m=1 & m=2 & m=1 & m=2 \\
\hline
0.1 & 1.779055 & - & 1.734042 & - \\
0.5 & 1.995825 & - & 2.189207 & - \\
1.5 & 2.630043 & - & 3.674960 & - \\
2.0 & 2.987921 & - & 4.530692 & - \\
\hline
\end{array}
$$

Figure 1. Graphs of the function $f(\eta)$ for $m=1, M=2$ and $S=0.1, 0.5, 1.0, 1.5, 2.0$ from bottom to top

Figure 2. Graphs of the function $f'(\eta)$ for $m=1, M=2$ and $S=0.1, 0.5, 1.0, 1.5, 2.0$ from bottom to top
Figure 3. Graphs of the function \( f(\eta) \) for \( m=1, S=1. \) and \( M=1, 2, 3 \) from bottom to top.

Figure 4. Graphs of the function \( f'(\eta) \) for \( m=1, S=1. \) and \( M=1, 2, 3 \) from bottom to top.

Figure 5. Graphs of the function \( f(\eta) \) for \( m=2, M=2 \) and \( S=1.5, 1, 1.5, 2 \) from bottom to top.

Figure 6. Graphs of the function \( f'(\eta) \) for \( m=2, M=2 \) and \( S=1, 1.5, 1, 1.5, 2 \) from bottom to top.

Figure 7. Graphs of the function \( f(\eta) \) for \( m=2, S=1. \) and \( M=1, 2, 3 \) from bottom to top.

Figure 8. Graphs of the function \( f'(\eta) \) for \( m=2, S=1. \) and \( M=1, 2, 3 \) from bottom to top.
REFERENCES


